

On an Initial Boundary Value Problem for a Generalized Degasperis-Procesi Equation

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Abstract: In this paper, we study the following initial boundary value problem for a generalized Degasperis-Procesi equation

$$\begin{cases} u_t - u_{xxt} + 4uu_x - 3u_x u_{xx} - uu_{xxx} = \gamma(u - u_{xx})_x, t \geq 0, x \in [0, 1] \\ u(0, t) = u(1, t) = u_x(0, t) = u_x(1, t) = 0, t \geq 0 \\ u(x, 0) = u_0(x), x \in [0, 1] \end{cases}$$

We establish local well-posedness to it by using Kato's theorem for abstract quasilinear evolution equation of hyperbolic type. By the energy estimate, we present a blow up result.

Key words: Degasperis-Procesi equation; initial boundary value problem; blow up

1 Introduction

Recently, Degasperis and Procesi [1] derived a new nonlinear dispersive partial differential equation

$$u_t - u_{xxt} + 4uu_x = 3u_x u_{xx} + uu_{xxx} \quad (1)$$

which is called Degasperis-Procesi equation. After Eq.(1) was derived, many papers were devoted to its study. For example, Degasperis, Holm and Hone [2] showed that the Eq.(1) is integrable by deriving a Lax pair and a bi-Hamiltonian structure for it. Dullin et al. [3] showed that the Degasperis-Procesi equation can be regarded as a model for nonlinear shallow water dynamics. Yin proved local well-posedness to Eq.(1) with initial data $u_0 \in H^S(R)$, $S > 3$ on the line [4] and on the circle [5]. The global existence of strong solutions and weak solutions to Eq.(1) are investigated in [5]-[11]. For the blow-up of solution to Cauchy problem for Eq.(1) we refer to [8]-[11]. Vakhnenko and Parkes [12] obtained the traveling wave solutions to Eq.(1). The shock wave solutions to Eq.(1) are investigated in [13]-[15]. Lenells [16] classified all weak traveling wave solutions. For more exact solutions to Eq.(1), we refer to [17]-[21]. For the results of the generalized Degasperis-Procesi equations, we refer to [22, 23, 25]. Escher and Yin [24] studied the initial boundary value problem of Eq.(1) on the half line R^+ and on finite interval $[0, 1]$. On $[0, 1]$, they used the boundary condition $u(0, t) = u(1, t) = 0$, and obtained the local existence of solution to Eq.(1). They also obtained that the solution $u(x, t)$ to Eq.(1) with boundary condition $u(0, t) = u(1, t) = 0$ and initial conditions $u_0 \in H^S(0, 1) \cap H_0^1(0, 1)$ ($\frac{3}{2} < S \leq 3$) blows up in finite time if $u_0 \neq 0$.

In this paper, we are concerned with the following initial boundary value problem for a generalized form of Eq.(1)

$$\begin{cases} u_t - u_{xxt} + 4uu_x - 3u_x u_{xx} - uu_{xxx} = \gamma(u - u_{xx})_x, t \geq 0, x \in [0, 1] \\ u(0, t) = u(1, t) = u_x(0, t) = u_x(1, t) = 0, t \geq 0 \\ u(x, 0) = u_0(x), x \in [0, 1] \end{cases} \quad (2)$$

where γ is a real constant, γu_x denotes the dissipative term and γu_{xxx} denotes the dispersive effect. Obviously, when $\gamma = 0$, Eq.(2) is the well known Degasperis-Procesi Eq.(1).

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The remainder of the paper is organized as follows. In Section 2, we establish the local well-posedness to the closed-loop system (2) by Kato’s theorem [26]. In Section 3, by using multiplier technique, we obtain the energy estimate. A blow-up phenomena is presented in this section.

We will use the following notation without further comment. $*$ for convolution; $L(Y, X)$ for all bounded linear operator from Banach space Y to X ($L(X)$ if $X = Y$); $\partial_x = \partial/\partial x$; $\Lambda = (1 - \partial_x^2)^{\frac{1}{2}}$; $H^S = H^S(0, 1)$ with the norm $\| \cdot \|_{H^S} = \| \cdot \|_S$; $L^2 = L^2(0, 1)$ with the norm $\| \cdot \|_{L^2}$; $H_{0,1}^S = \{u \in H^S(0, 1) : u(0, t) = u(1, t)\}$; $[A, B]$ denotes the commutator of the linear operators A and B ; $C^k(I; X)$ for the space of all k times continuously differentiable functions defined on an interval I with values in Banach space X .

2 Local well-posedness

In this section, we will apply Kato’s theorem [26] to establish the local well-posedness for the closed-loop system (2). For convenience, we state Kato’s theorem in the form suitable for our purpose.

Consider the Cauchy problem associated to a quasilinear evolution equation

$$\begin{cases} \frac{du}{dt} + A(u)u = f(u) \in X, & t \geq 0 \\ u(0) = u_0 \in Y, \end{cases} \tag{1}$$

where $A(u)$ is a linear operator depending on the unknown u , and u_0 the initial value. To study the Cauchy problem (local in the time) associated to (1), we will make the following assumptions:

(X) X and Y are reflexive Banach spaces where $Y \subset X$, with the inclusion continuous and dense, and there is an isomorphism Q from Y onto X such that $\|\phi\|_Y = \|Q\phi\|_X$ for all $\phi \in Y$.

(A₁) Let W be an open ball centered in 0 and contained in Y . The linear operator $A(u)$ belongs to $G(X, 1, \beta)$ where β is a real number, i.e., $-A(u)$ generates a C_0 -semigroup such that $\|e^{-sA(u)}\|_{B(X)} \leq e^{\beta s}$.

Note that if X is a Hilbert space, then $A \in G(X, 1, \beta)$ if and only if

- (a) $(A\phi, \phi)_X \geq -\beta \|\phi\|_X^2, \forall \phi \in D(A)$,
- (b) $(A + \lambda I)$ is onto for some (all) $\lambda > \beta$.

Under these conditions $A(u)$ is said to be quasi-m-accretive.

(A₂) The map $w \in W \rightarrow B(w) = [Q, A(w)]Q^{-1} \in L(X)$ is uniformly bounded and Lipschitz continuous, that is, there exist constants $\lambda_1, \mu_1 > 0$, such that for all $w, y \in W$,

$$\|B(w)\|_{L(X)} \leq \lambda_1,$$

(A₃) $X \subseteq D(A(w))$ for each $w \in W$ (so that $A(w)|_Y \in L(Y, X)$ by the Closed Graph theorem). Moreover, the map $w \in W \rightarrow A(w) \in L(Y, X)$ satisfies the following Lipschitz condition:

$$\|A(w) - A(y)\|_{L(Y, X)} \leq \mu_2 \|w - y\|_X,$$

for all $w, y \in W$, where μ_2 is a non-negative constant.

(f) The function $f : W \rightarrow Y$ is bounded, i.e., there is a constant $\lambda_2 > 0$ such that $\|f(w)\|_Y \leq \lambda_2$ for all $w \in W$, and the function $w \in X \rightarrow f(w)$ is Lipschitz in X (resp. in Y), i.e.

$$\begin{aligned} \|f(w) - f(y)\|_X &\leq \mu_3 \|w - y\|_X, \quad \forall w, y \in W, \\ \|f(w) - f(y)\|_Y &\leq \mu_4 \|w - y\|_Y, \quad \forall w, y \in W, \end{aligned}$$

where μ_3, μ_4 is non-negative constant.

We are now in position to state Kato’s local well posedness result.

Theorem 1 (Kato’s theorem [26]) *Assume conditions (X), (A₁)-(A₃) and (f) hold. Given $u_0 \in Y$, there is $T > 0$ and unique solution $u \in C([0, T]; Y) \cap C^1([0, T]; X)$ to (1) with $u(0) = u_0$. Moreover, the map $u_0 \in Y \rightarrow u \in C([0, T]; Y)$ is continuous.*

We now provide the framework in which we shall reformulate problem (2).

Let $m = u - u_{xx}$, then the first equation of (2) takes the form of a quasi-linear evolution equation of hyperbolic type

$$m_t + um_x + 3u_x m = \gamma m_x. \tag{2}$$

By using the operator $G(x) = \frac{\cosh(x-[x]-\frac{1}{2})}{2 \sinh(\frac{1}{2})}$, where $[x]$ denotes the integer part of $x \in [0, 1]$, then $(1 - \partial_x^2)^{-1} f = G * f, \forall f \in L^2$ and $G * m = u$. Using this identity, we can rewrite Eq.(2) as

$$u_t + uu_x + \partial_x(G * (\frac{3}{2}u^2)) = \gamma u_x$$

Then the closed-loop system (2) becomes

$$\begin{cases} u_t + uu_x + \partial_x(G * (\frac{3}{2}u^2)) = \gamma u_x, & t \geq 0, x \in [0, 1] \\ u(0, t) = u(1, t) = u_x(0, t) = u_x(1, t) = 0, & t \geq 0 \\ u(x, 0) = u_0(x), & x \in [0, 1] \end{cases} \tag{3}$$

Theorem 2 Given $u_0(x) \in H^2_{0,1}$, there exists a maximal time value $T = T(u_0(x)) > 0$ and a unique solution $u(x, t)$ to (3) such that $u = u(\cdot ; u_0) \in C([0, T]; H^2_{0,1}) \cap C^1([0, T]; H^1)$. Moreover, the solution depends continuously on the initial data, i.e., the mapping $u_0 \rightarrow u(\cdot ; u_0) : H^2_{0,1} \rightarrow C([0, T]; H^2_{0,1}) \cap C^1([0, T]; H^1)$ is continuous. Equivalently, system (2) has a unique solution $u(x, t) \in C([0, T]; H^2_{0,1}) \cap C^1([0, T]; H^1)$.

Proof. Define $A(u)u = u\partial_x u - \gamma\partial_x u, f(u) = -\partial_x(G * (\frac{3}{2}u^2)), Q = \Lambda, X = H^1, Y = H^2_{0,1}$. Then the conditions (X), (A_1) - (A_3) and (f) in Theorem 2.1 are satisfied. Here we omit the details. One can see the similar argument in [26]. ■

3 blow up

By using multiplier technique, we obtain the following energy estimate of the closed-loop system (3).

Theorem 3 Assume $u_0(x) \in H^2_{0,1}$. Let T be the maximal existence time of the solution $u(x, t)$ to (3) guaranteed by Theorem 2.2. Then we have $\|u(\cdot, t)\|_{L^2}^2 \leq 4\|u_0\|_{L^2}^2$.

Proof. For $u \in H^2_{0,1}$, we define the function v as follow

$$4v - v_{xx} = u.$$

Let $G_1(x) = e^{-2|x|}$ be the Green's function of the operator $4 - \partial_x^2$, then

$$u(x, t) = (G_1 * v)(x, t) = \int_0^1 e^{-2|x-y|} v(y, t) dy$$

Multiplying the first equation of (3) by $v - v_{xx}$, and integrating over $(0, 1)$ by parts, we get

$$\begin{aligned} \int_0^1 u_t(v - v_{xx})dx &= - \int_0^1 uu_x(v - v_{xx})dx - \int_0^1 \partial_x(G * (\frac{3}{2}u^2))(v - v_{xx})dx \\ &\quad + \gamma \int_0^1 u_x(v - v_{xx})dx \end{aligned} \tag{1}$$

For the LHS of (1), by the periodic boundary conditions, using (1) and integration by parts, we have

$$\begin{aligned} \int_0^1 u_t(v - v_{xx})dx &= \int_0^1 (4v_t - v_{txx})(v - v_{xx})dx \\ &= -4v_t v \Big|_0^1 - v_{tx} v \Big|_0^1 + \int_0^1 (4v_t v + 5v_{tx} v_x + v_{txx} v_{xx})dx \\ &= \frac{1}{2} \frac{d}{dt} \int_0^1 (4v^2 + 5v_x^2 + v_{xx}^2)dx \end{aligned} \tag{2}$$

Similarly, for the RHS of (1), we obtain

$$\begin{aligned}
 & - \int_0^1 uu_x(v - v_{xx})dx - \int_0^1 \partial_x(G * (\frac{3}{2}u^2))(v - v_{xx})dx + \gamma \int_0^1 u_x(v - v_{xx})dx \\
 & = - \int_0^1 uu_x(v - v_{xx})dx - vG * (\frac{3}{2}u^2) |_0^1 + \int_0^1 (G * (\frac{3}{2}u^2))v_x dx \\
 & + \partial_x(G * (\frac{3}{2}u^2))v_x |_0^1 - \int_0^1 \partial_x^2(G * (\frac{3}{2}u^2))v_x dx + \gamma \int_0^1 (4v_x - v_{xx})(v - v_{xx})dx \\
 & = - \int_0^1 uu_x(v - v_{xx})dx + \int_0^1 \frac{3}{2}u^2v_x dx + 2\gamma v^2|_0^1 - 2\gamma v_x^2|_0^1 - \gamma vv_{xx}|_0^1 + \frac{\gamma}{2}v_x^2|_0^1 + \frac{\gamma}{2}v_{xx}^2|_0^1 \\
 & = - \int_0^1 uu_x(v - v_{xx})dx + \frac{3}{2}u^2v |_0^1 - 3 \int_0^1 uu_xv dx \\
 & = - \int_0^1 uu_x(4v - v_{xx})dx \\
 & = - \int_0^1 u^2u_x dx \\
 & = -\frac{1}{3}u^3 |_0^1 = 0
 \end{aligned} \tag{3}$$

It follows from (1)-(3) that

$$\frac{1}{2} \frac{d}{dt} \int_0^1 (4v^2 + 5v_x^2 + v_{xx}^2)dx = 0,$$

i.e.,

$$\begin{aligned}
 & 4 \|v(\cdot, t)\|_{L^2}^2 + 5 \|v_x(\cdot, t)\|_{L^2}^2 + \|v_{xx}(\cdot, t)\|_{L^2}^2 \\
 & = 4 \|v(\cdot, 0)\|_{L^2}^2 + 5 \|v_x(\cdot, 0)\|_{L^2}^2 + \|v_{xx}(\cdot, 0)\|_{L^2}^2
 \end{aligned} \tag{4}$$

On the other hand, we have

$$\begin{aligned}
 \|u(\cdot, t)\|_{L^2}^2 & = \int_0^1 u^2 dx = \int_0^1 (4v - v_{xx})^2 dx \\
 & = 16 \int_0^1 v^2 dx - 8vv_x |_0^1 + 8 \int_0^1 v_x^2 dx + \int_0^1 v_{xx}^2 dx \\
 & = 16 \|v(\cdot, t)\|_{L^2}^2 + 8 \|v_x(\cdot, t)\|_{L^2}^2 + \|v_{xx}(\cdot, t)\|_{L^2}^2 \\
 & \geq 4 \|v(\cdot, t)\|_{L^2}^2 + 5 \|v_x(\cdot, t)\|_{L^2}^2 + \|v_{xx}(\cdot, t)\|_{L^2}^2.
 \end{aligned} \tag{5}$$

Then it follows from (4) and (5) that

$$\begin{aligned}
 \|u(\cdot, t)\|_{L^2}^2 & = \int_0^1 u^2 dx = 16 \|v(\cdot, t)\|_{L^2}^2 + 8 \|v_x(\cdot, t)\|_{L^2}^2 + \|v_{xx}(\cdot, t)\|_{L^2}^2 \\
 & \leq 4(4 \|v(\cdot, t)\|_{L^2}^2 + 5 \|v_x(\cdot, t)\|_{L^2}^2 + \|v_{xx}(\cdot, t)\|_{L^2}^2) \\
 & = 4(4 \|v(\cdot, 0)\|_{L^2}^2 + 5 \|v_x(\cdot, 0)\|_{L^2}^2 + \|v_{xx}(\cdot, 0)\|_{L^2}^2) \\
 & \leq 4 \|u_0\|_{L^2}^2.
 \end{aligned} \tag{6}$$

■

Now we present a blow-up result of solution to the closed-loop system (2) (or system (3)).

Theorem 4 Assume $u_0(x) \in H_{0,1}^2$. Let T be the maximal existence time of the solution $u(x, t)$ to (3) guaranteed by Theorem 2.2. If there exists one point $x_0 \in (0, 1)$ such that $u_{xx}(x_0, t) = 0$ and $v'_0(x_0) < -\sqrt{6}\|u_0\|_{L^2}$, then the corresponding solution to (3) blows up in finite time. Moreover, the maximal time of existence is estimated above by

$$\frac{1}{\sqrt{6}\|u_0\|_{L^2}} \ln\left(\frac{h(0) - \sqrt{6}\|u_0\|_{L^2}}{h(0) + \sqrt{6}\|u_0\|_{L^2}}\right),$$

where $h(0) = u'_0(x_0)$.

Proof. Let $T > 0$ be the maximal time of existence of the solution $u(x, t)$ to (3) with the initial data $u_0(x) \in H^2_{0,1}$. Differentiating the first equation of (3) with respect to x , in view of $\partial_x^2 G * f = G * f - f$, we have

$$u_{tx} = -u_x^2 - uu_{xx} + \gamma u_{xx} + \frac{3}{2}u^2 - G * (\frac{3}{2}u^2). \tag{7}$$

Let $x = x_0$ in (7) and set $h(t) = u_x(x_0, t)$. Noting that $G * (\frac{3}{2}u^2) \geq 0$, we obtain

$$h'(t) = -h^2(t) + \frac{3}{2}u^2(x_0, t). \tag{8}$$

In view of (6), we have

$$u^2(x_0, t) \leq \|u\|_{L^\infty}^2 \leq \|u\|_{L^2}^2 \leq 4\|u_0\|_{L^2}^2. \tag{9}$$

It follows from (8) and (9) that

$$h'(t) \leq -h^2(t) + 6\|u_0\|_{L^2}^2.$$

Note that if $h(0) \leq -\sqrt{6}\|u_0\|_{L^2}$, then $h(t) \leq \sqrt{6}\|u_0\|_{L^2}$, for all $t \in [0, T)$. Therefore, from the above inequality we obtain

$$\frac{h(0) + \sqrt{6}\|u_0\|_{L^2}}{h(0) - \sqrt{6}\|u_0\|_{L^2}} e^{\sqrt{6}\|u_0\|_{L^2} t} - 1 \leq \frac{\sqrt{6}\|u_0\|_{L^2}}{h(t) - \sqrt{6}\|u_0\|_{L^2}} < 0.$$

Due to $0 < \frac{h(0) + \sqrt{6}\|u_0\|_{L^2}}{h(0) - \sqrt{6}\|u_0\|_{L^2}} < 1$, then exists

$$T_0 \leq \frac{1}{\sqrt{6}\|u_0\|_{L^2}} \ln\left(\frac{h(0) - \sqrt{6}\|u_0\|_{L^2}}{h(0) + \sqrt{6}\|u_0\|_{L^2}}\right)$$

such that

$$\lim_{t \rightarrow T_0} h(t) = -\infty.$$

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